Reference is also made to the treatise by Luke [2], which includes a discussion of this integral, and gives an equivalent expression in closed form when $b=1$.
J. W. W.

1. Norio Kato, "Integrated intensities of the diffracted and transmitted X-rays due to ideally perfect crystals (Laue case)," J. Phys. Soc. Japan, v. 10, 1955, p. 46-55.
2. Yudell L. Luke, Integrals of Bessel Functions, McGraw-Hill Book Co., 1962, p. 122 and Chapter X.

15[L, M].-N. Skoblia, Tables for the Numerical Inversion of Laplace Transforms, Academy of Sciences of USSR, Moscow, 1964, 44 p., 22 cm . Paperback. Price 13 kopecks.
Consider the Laplace transform pair (which we assume exists)

$$
\begin{equation*}
p^{-s} g(p)=\int_{0}^{\infty} e^{-p t} f(t) d t, \quad f(t)=(2 \pi i)^{-1} \int_{c-i \infty}^{c+i \infty} p^{-s} e^{p t} g(p) d p \tag{1}
\end{equation*}
$$

where $c>0$ and $c$ lies to the right of all singularities of $g(p)$. Suppose that $g(p)$ is known and can be represented by a polynomial in $1 / p$. Then an approximation ${ }^{\prime}$ formula for $f(t)$ is readily constructed from the second formula in (1). Now it may be shown that

$$
\begin{equation*}
\int_{c-i \infty}^{c+i \infty} e^{p} p^{-s} P_{n}\left(p^{-1}\right) P_{m}\left(p^{-1}\right) d p=\delta_{m n} h_{n} \tag{2}
\end{equation*}
$$

$$
h_{n}=\frac{(-1)^{n} n!}{(2 n-1+s) \Gamma(n-1+s)}
$$

where $\delta_{m n}$ is the Kronecker delta, and in hypergeometric notation,

$$
\begin{align*}
P_{n}(x) & =2 F_{0}(-n, n+s-1 ; x) \\
& =(2-2 n-s)_{n} x^{n}{ }_{1} F\left(-n ; 2-2 n-s ; \frac{1}{x}\right) \tag{3}
\end{align*}
$$

This shows that numerous properties of $P_{n}(x)$ follow from known results on confluent hypergeometric functions. In view of (2), we have the approximation

$$
\begin{equation*}
f(t) \sim(2 \pi i)^{-1} \sum_{k=1}^{n} A_{k, n} g\left(p_{k}\right) \tag{4}
\end{equation*}
$$

where

$$
\begin{equation*}
P_{n}\left(p_{k}{ }^{-1}\right)=0, \quad k=1,2, \cdots, n \tag{5}
\end{equation*}
$$

and the weights, $A_{k, n}$ are the Christoffel numbers. Thus, the approximation is exact if indeed $g(p)$ is a polynomial in $1 / p$ of degree $(2 n-1)$. A convenient formula for the weights is

$$
\begin{equation*}
A_{k, n}=\sum_{m=0}^{n-1}\left\{P_{m}\left(p_{n}{ }^{-1}\right)\right\}^{2} / h_{m} \tag{6}
\end{equation*}
$$

The pamphlet gives some properties of $P_{n}(x)$, though (3) and (6) are not among them. The following are tabulated to $7 \mathrm{~S}: p_{k}, A_{k, n}$ for $k=1(1) n, n=1(1) 10$, and $s=0.1(0.1) 3.0$.

The case $s=1$ has been treated by H. Salzer. (See "Orthogonal polynomials arising in the numerical evaluation of Laplace transforms," MTAC, v. 9, 1955, p. 164-177, and "Additional formulas and tables for orthogonal polynomials originating from inversion integrals," J. Math. Phys., v. 40, 1961, p. 72-86.) These latter sources give the zeros and weights to 15 D for $n=1(1) 15$. Note that Salzer's quadrature formula is exact if $g(p)$ is a polynomial in $1 / p$ of degree $2 n$ such that $g(\infty)=0$. In the booklet under review, the quadrature formula is exact if $g(p)$ is of degree $(2 n-1)$, but $g(\infty)$ need not vanish. Thus the Christoffel numbers in Salzer's work differ from those of the present author. However, the zeros are the same. Twice the negatives of the zeros of $P_{n}(x)$ have been tabulated mostly to 5 D by V. N. Kublanovskaia and T. N. Smirnova. (See "Zeros of Hankel functions and some related functions," Trudy. Mat. Inst. AN, USSR No. 53, 1959, p. 186192. This is also available as Electronic Research Directorate, Air Force Cambridge Research Laboratories Report AFCRL-TN 60-1128, October 1960.)
Y. L. L.

16[M, X].-E. L. Albasiny, R. J. Bell \& J. R. A. Cooper, A Table for the Evaluation of Slater Coefficients and Integrals of Triple Products of Spherical Harmonics, National Physical Laboratory Mathematics Division Report No. 49, 1963, $x i+163$ pages.
This is a table of integrals

$$
\int_{-1}^{+1} \Theta_{l_{1}}^{m_{1}}(x) \Theta_{l_{2}}^{m_{2}}(x) \Theta_{l_{3}}^{m_{3}}(x) d x
$$

where
$\Theta_{l}{ }^{m}(x)=(-1)^{m}\left[\frac{(2 l+1)}{2} \frac{(l-m)!}{(l+m)!}\right]^{1 / 2} \frac{\left(1-x^{2}\right)^{m / 2}}{l!2^{l}} \frac{d^{l+m}}{d x^{l+m}}\left(x^{2}-1\right)^{l}$,

$$
0 \leqq m \leqq l
$$

is an associated Lengendre function. These integrals are closely related to integrals which occur in molecular structure calculations (see, for example, [1]).

Values of the above integrals are tabulated to 12 decimal places for ( $m_{i}, l_{i}$ integers)

$$
\begin{gathered}
m_{1} \pm m_{2} \pm m_{3}=0 \\
l_{1} \leqq l_{2} \leqq l_{3} \\
l_{1}+l_{2}+l_{3} \text { even } \\
\left|l_{1}-l_{2}\right| \leqq l_{3} \leqq l_{1}+l_{2} \\
l_{1}, l_{2} \leqq 12, l_{3} \leqq 24
\end{gathered}
$$

Under these conditions the integrand is a polynomial of degree $\leqq 48$, and thus can be calculated exactly, using an $n$-point Gauss-Legendre quadrature formula [3, p. $107-111]$ for $n \geqq 25$. The tables were computed using the 25 -point formula tabulated by Gawlik [2] and recomputed as a check using the 26-point formula. The calculations were carried out on the ACE computer, which has a 46 -bit floating-

